

Optimal Design of Tube Banks in Crossflow Using Entropy Generation Minimization Method

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An entropy generation minimization, EGM, method is applied as a unique measure to study the thermodynamic losses caused by heat transfer and pressure drop for a fluid in crossflow with tube banks. The use of EGM allows the combined effect of heat transfer and pressure drop to be assessed through the simultaneous interaction with the tube bank. A general dimensionless expression for the entropy generation rate is obtained by considering a control volume around a tube bank and applying the conservation equations for mass and energy with the entropy balance. Analytical/empirical correlations for heat transfer coefficients and friction factors are used, where the characteristic length is used as the diameter of the tubes and reference velocity used in Reynolds number and pressure drop is based on the minimum free area available for the fluid flow. Both in-line and staggered arrangements are studied and their relative performance is compared for the same thermal and hydraulic conditions. A parametric study is also performed to show the effects of different design variables on the overall performance of tube banks. It is shown that all relevant design parameters for tube banks, including geometric parameters and flow conditions can be simultaneously optimized.

Nomenclature

A	surface area of a single tube, m^2
A_t	total heat transfer area, m^2
C_1	constant defined in Eqs. (15)
D	tube diameter, m
e	specific energy, W
f	friction factor
g, l	equality and inequality constraints
h_{avg}	average heat transfer coefficient of tubes, $W/m^2 \cdot K$
j	number of imposed constraints
K_1	constant defined in Eqs. (17)
k	thermal conductivity, W/mK
\mathcal{L}	Lagrangian function
L	length of tube, m
N	total number of tubes $\equiv N_T N_L$

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n	number of design variables
N_L	number of rows in streamwise direction
N_s	dimensionless entropy generation rate $\equiv \dot{S}_{gen}/(Q^2 U_{max}/k_f \nu T_a^2)$
N_T	number of rows in spanwise direction
Nu_D	Nusselt number based on tube diameter
P	pressure, Pa
\dot{Q}	heat transfer rate over the boundaries of CV, W
Re_D	Reynolds number $\equiv DU_{max}/\nu$
\dot{S}_{gen}	total entropy generation rate, W/K
S_D	dimensionless diagonal pitch $\equiv S_D/D$
S_L	dimensionless streamwise pitch $\equiv S_L/D$
S_T	dimensionless spanwise pitch $\equiv S_T/D$
S_D	diagonal pitch, m
S_L	tube spacing in streamwise direction, m
S_T	tube spacing in spanwise direction, m
T	absolute temperature, K
U_{app}	approach velocity, m/s
U_{max}	maximum velocity in minimum flow area, m/s
Pr	Prandtl number $\equiv \nu/\alpha$
Re_D	Reynolds number based on maximum velocity in minimum flow area $\equiv DU_{max}/\nu$
S_D	diagonal pitch $[m]$
S_L	longitudinal distance between two consecutive tubes, m
S_T	transverse distance between two consecutive tubes, m
S_L	dimensionless longitudinal pitch $\equiv S_L/D$
S_T	dimensionless transverse pitch $\equiv S_T/D$
v	specific volume of fluid, m^3/kg
\dot{W}_{cv}	energy transfer by work across the boundaries of CV, W
x_i	design variables

Greek Symbols

γ	aspect ratio $\equiv L/D$
ν	kinematic viscosity of fluid $[m^2/s]$
ρ	fluid density $[kg/m^3]$

Subscripts

a	ambient
f	fluid
in	inlet of CV
out	exit of CV
T	thermal
w	wall

I. Introduction

DUE to extensive use of high performance compact heat exchangers, an optimal design of tube banks is very important. These heat exchangers are found in numerous applications, such as an automobile radiator, an oil cooler, a preheater, an air-cooled steam condenser, a shell and tube type heat exchanger, and the evaporator of an air-conditioning system. Tube banks are usually arranged in an in-line or staggered

manner, where one fluid moves across the tubes, and the other fluid at a different temperature passes through the tubes. In this study, the authors are specifically interested to determine an optimal design of the tube banks in crossflow using entropy generation minimization method. The crossflow correlations for the heat transfer and pressure drop are employed to calculate entropy generation rate.

A. Literature Review

A careful review of existing literature reveals that most of the studies are related to the optimization of plate heat exchangers and only few studies are related to tube heat exchangers.

To the authors knowledge, McClintock¹ was the first one who employed the concept of irreversibility for estimating and minimizing the usable energy wasted in heat exchangers design. Later on, Bejan^{2,3,4,5} extended that concept and presented an optimum design method for balanced and imbalanced counterflow heat exchangers. He proposed the use of a “Number of Entropy Production Units” N_s as a basic parameter in describing heat exchanger performance. This method was applied to a shell and tube regenerative heat exchanger to obtain the minimum heat transfer area when the amount of units was fixed. Later on, Aceves-Saborio et al.⁶ extended that approach to include a term to account for the exergy of the heat exchanger material. Grazzini and Gori⁷, Sekulic⁸, Zhang et al.⁹, Ordóñez and Bejan¹⁰, Bejan^{11,12} demonstrated that the optimal geometry of a counterflow heat exchanger can be determined based on thermodynamic optimization subject to volume constraint. Yilmaz et al.¹³ first recalled and discussed the need for the systematic design of heat exchangers using a second law-based procedure and then presented second-law based performance evaluation criteria for heat exchangers. Entropy generation rate is generally used in a dimensionless form. Unuvar and Kargici¹⁴ and Peters and Timmerhaus¹⁵ presented an approach for the optimum design of heat exchangers. They used the method of steepest descent for the minimization of annual total cost. They observed that this approach is more efficient and effective to solve the design problem of heat exchangers. Optimization of plate-fin and tube-fin crossflow heat exchangers was presented by Shah et al.^{16,17}, Kröger¹⁸, and Van den Bulck¹⁸. They employed optimal distribution of the UA value across the volume of crossflow heat exchangers and optimized different design variables like fin thickness, fin height, and fin pitch.

Cylinder geometry was optimized in a paper by Poulikakos and Bejan¹⁹ and Khan et al.²⁰. After the general formula was derived using the EGM method analytical methods and graphical results were developed which resulted in the optimum selection of the dimensions of several different fin configurations. Bejan et al.²¹ showed that EGM may be used by itself in the preliminary stages of design, in order to identify trends and the existence of optimization opportunities. In two different studies, Stanescu et al.²², and Matos et al.²³ demonstrated that the geometric arrangement of tubes/cylinders in cross-flow forced convection can be optimized for maximum heat transfer subject to overall volume constraint. They used FEM to show the optimal spacings between rows of tubes. Vargas, et al.²⁴ documented the process of determining the internal geometric configuration of a tube bank by optimizing the global performance of the installation that uses the crossflow heat exchanger.

B. Assumptions

The following assumptions are used in this study:

1. Tube bank is insulated from its surroundings.
2. Tubes are plain.
3. Flow is 2-D, steady, and laminar.
4. Fluid is Newtonian and incompressible with constant properties.
5. Conduction along tube walls is negligible.
6. Radiation heat transfer from tubes is negligible.
7. Potential and kinetic energy changes are negligible.

II. Model Development

Consider the in-line tube bank in cross-flow delineated by a control volume (CV) in Figure 1. The sides of this control volume can be regarded as impermeable, adiabatic and shear free. Fluid I moves across the tubes while fluid II at a different temperature passes through the tubes. In this study, the authors are concerned only with fluid I. The approach velocity of the fluid I is U_{app} and the ambient temperature is T_a . The temperature of the tube wall is T_w . The properties of the fluid are represented by e_{in} , P_{in} , s_{in} at the inlet and by e_{out} , P_{out} , s_{out} at the outlet respectively. The irreversibility of this system is also due to heat transfer across the nonzero temperature difference $T_w - T_a$ and due to the total pressure drop across the tube bank. The mass rate balance for the CV, shown in Fig. 1, gives

$$\frac{dm_{cv}}{dt} = \dot{m}_{in} - \dot{m}_{out} \quad (1)$$

For steady state, it reduces to

$$\dot{m}_{in} = \dot{m}_{out} = \dot{m} \quad (2)$$

First law of thermodynamics for the same CV can be written as

$$\frac{dE_{cv}}{dt} = \dot{Q} - \dot{W}_{cv} + \dot{m}_{in} (e_{in} + P_{in}v_{in}) - \dot{m}_{out} (e_{out} + P_{out}v_{out}) \quad (3)$$

where $\frac{dE_{cv}}{dt}$ is the time rate of change of energy

within the CV and for steady state, $\frac{dE_{cv}}{dt} = 0$. The specific energy e is the sum of specific internal, kinetic, and potential energies. Due to continuity and same elevation of the CV, $V_{in} = V_{out}$ and $z_{in} = z_{out}$, so the kinetic and potential energy terms will drop out. Therefore, $e_{in} = u_{in}$ and $e_{out} = u_{out}$. The only work is flow work at the inlet and exit of the CV, so the term \dot{W}_{cv} also drops out. Thus the energy rate balance reduces to:

$$\dot{Q} = \dot{m} \left[\underbrace{(u_{out} + P_{out}v_{out})}_{h_{out}} - \underbrace{(u_{in} + P_{in}v_{in})}_{h_{in}} \right] \quad (4)$$

The combination of specific internal and flow energies is defined as specific enthalpy; therefore, the energy rate balance reduces further to:

$$\dot{Q} = \dot{m}(h_{out} - h_{in}) \quad (5)$$

From the second law of thermodynamics

$$\frac{dS_{cv}}{dt} = \dot{m}(s_{in} - s_{out}) + \frac{\dot{Q}}{T_w} + \dot{S}_{gen} \quad (6)$$

For steady state, $\frac{dS_{cv}}{dt} = 0$, so the entropy rate balance reduces to

$$\dot{S}_{gen} = \dot{m}(s_{out} - s_{in}) - \frac{\dot{Q}}{T_w} \quad (7)$$

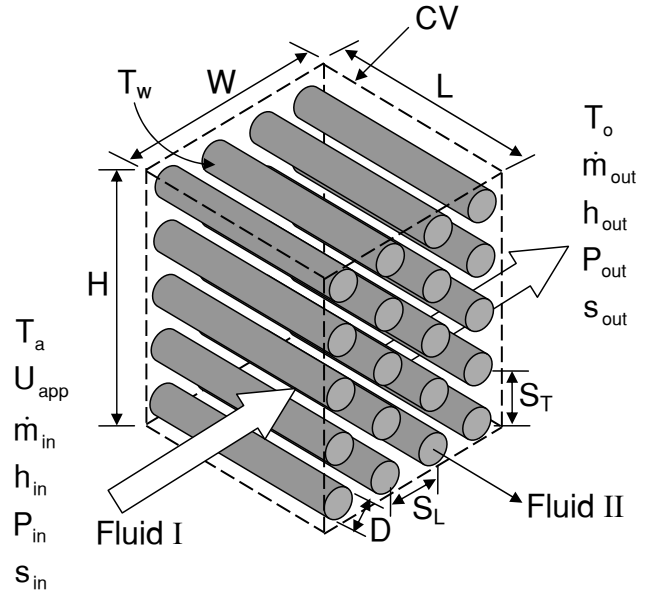


Figure 1. Control Volume for Calculating \dot{S}_{gen} for Tube Banks.

Gibb's equation $[dh = Tds + (1/\rho)dP]$ can be written as:

$$h_{out} - h_{in} = T_a(s_{out} - s_{in}) + \frac{1}{\rho}(P_{out} - P_{in}) \quad (8)$$

Combining Eqs. (5) and (8), we get:

$$\mathcal{Q} = \dot{m}T_a(s_{out} - s_{in}) - \frac{\dot{m}\Delta P}{\rho} \quad (9)$$

Combining Eqs. (7) and (9), we get:

$$\dot{S}_{gen} = \left(\frac{\mathcal{Q}^2}{T_a T_w} \right) R_{tube} + \frac{\dot{m}\Delta P}{\rho T_a} \quad (10)$$

where R_{tube} is the tube wall thermal resistance, \dot{m} is the mass flow rate through the tubes and ΔP is the pressure drop across the tube bank and can be written as

$$R_{tube} = \frac{\Delta T}{\mathcal{Q}} = \frac{1}{h_{avg}A} \quad (11)$$

$$\dot{m} = \rho U_{app} N_T S_T L \quad (12)$$

$$\Delta P = N_L f \left(\frac{1}{2} \rho U_{max}^2 \right) \quad (13)$$

Khan [23] has developed following analytical correlation for dimensionless heat transfer coefficient for the tube banks:

$$Nu_D = \frac{h_{avg}D}{k_f} = C_1 Re_D^{1/2} Pr^{1/3} \quad (14)$$

where Re_D is the Reynolds number defined as $Re_D = \frac{DU_{max}}{\nu}$ and C_1 is a constant which depends upon the longitudinal and transverse pitch ratios, arrangement of the tubes, and thermal boundary conditions. For isothermal boundary condition, it is given by:

$$C_1 = \begin{cases} [0.25 + \exp(-0.55\mathcal{S}_L)]\mathcal{S}_T^{0.285}\mathcal{S}_L^{0.212} & \text{In-Line} \\ \frac{0.61\mathcal{S}_T^{0.091}\mathcal{S}_L^{0.053}}{[1 - 2\exp(-1.09\mathcal{S}_L)]} & \text{Staggered} \end{cases} \quad (15)$$

Equation (15) is valid for $1.05 \leq \mathcal{S}_L \leq 3$ and $1.05 \leq \mathcal{S}_T \leq 3$.

Žukauskas and Ulinskas¹⁸ collected data, from a variety of sources, about friction factors for flow in the in-line and staggered arrangements having many rows and plotted them in the form Eu/K_1 versus Re_D , where K_1 is a correction factor accounting for geometry. They fitted these plots by inverse power series relationships and recommended more than 30 correlations for friction and correction factors depending on the transverse and longitudinal pitch ratios, the Reynolds number range and the type of arrangement. Khan¹ digitized thier experimental data and fitted into single correlations for the friction and correction factors for each arrangement. These correlations can be used for any pitch ratio $1.05 \leq \mathcal{S}_L$ or $\mathcal{S}_T \leq 3.0$ and Reynolds number in the laminar flow range. They are

$$f = \begin{cases} K_1 [0.233 + 45.78/(\mathcal{S}_T - 1)^{1.1} Re_D] & \text{In-Line} \\ K_1 [378.6/\mathcal{S}_T^{13.1/\mathcal{S}_T}] / Re_D^{0.68/\mathcal{S}_T^{1.29}} & \text{Staggered} \end{cases} \quad (16)$$

where K_1 is a correction factor depending upon the flow geometry and arrangement of the pins. It is given by:

$$K_1 = \begin{cases} 1.009 \left(\frac{\mathcal{S}_T - 1}{\mathcal{S}_L - 1} \right)^{1.09/Re_D^{0.0553}} & \text{In-Line} \\ 1.175(\mathcal{S}_L/\mathcal{S}_T Re_D^{0.3124}) + 0.5 Re_D^{0.0807} & \text{Staggered} \end{cases} \quad (17)$$

The velocity U_{max} , used in Eq. (13) and in the definition of Reynolds number, represents the maximum average velocity seen by the bank as flow accelerates between tubes, and is given by:

$$U_{max} = \max \left\{ \frac{\mathcal{S}_T}{\mathcal{S}_T - 1} U_{app}, \frac{\mathcal{S}_T}{\mathcal{S}_D - 1} U_{app} \right\} \quad (18)$$

where $\mathcal{S}_D = \sqrt{\mathcal{S}_L^2 + (\mathcal{S}_T/2)^2}$ is the dimensionless diagonal pitch.

Using Eqs. (11) - (14), the entropy generation rate can be simplified to

$$\dot{S}_{gen} = \frac{\mathcal{Q}^2/T_a T_w}{C_1 N \pi L k_f Re_D^{1/2} Pr^{1/3}} + \frac{N f \rho U_{max}^3 (\mathcal{S}_T - 1) L}{2 T_a} \quad (19)$$

The first term on the RHS represents the entropy generation rate due to heat transfer whereas the second term represents the entropy generation rate due to fluid flow. For external flow, Bejan (1996) used the term $\mathcal{Q}^2 U_{max}/k_f \nu T_a^2$ to nondimensionalize entropy generation rate in Eq. (19). So the dimensionless entropy generation rate can be written as

$$N_s = \frac{T_a/T_w}{C_1 N \pi \gamma Re_D^{3/2} Pr^{1/3}} + \frac{1}{2} f N \gamma B Re_D^2 (\mathcal{S}_T - 1) \quad (20)$$

where $B = \rho \nu^3 k_f T_a / \mathcal{Q}^2$ is a fixed dimensionless duty parameter that accounts for the importance of fluid friction irreversibility relative to heat transfer irreversibility. Equation (20) shows that, for the given volume of the tube bank and heat duty, the dimensionless entropy generation rate depends on ambient and wall temperatures, total number of tubes, longitudinal and transverse pitch ratios, Reynolds and Prandtl numbers, and aspect ratio. After fixing ambient and wall temperatures, all these parameters depend on tube diameter and the approach velocity for given longitudinal and transverse pitches.

III. Optimization Procedure

The problem considered in this study is to minimize the dimensionless entropy generation rate, given by Eq. (20), for the optimal overall performance of the tube bank. If $f(\mathbf{x})$ represent the dimensionless entropy generation rate that is to be minimized subject to equality constraints $g_j(x_1, x_2, \dots, x_n) = 0$ and inequality constraints $l_k(x_1, x_2, \dots, x_n) \geq 0$, then the complete mathematical formulation of the optimization problem may be written in the following form:

$$\text{minimize } f(\mathbf{x}) = N_s(\mathbf{x}) \quad (21)$$

subject to the equality constraints

$$g_j(\mathbf{x}) = 0, \quad j = 1, 2, \dots, m \quad (22)$$

and inequality constraints

$$l_j(\mathbf{x}) \geq 0, \quad j = m + 1, \dots, n \quad (23)$$

where g_j and l_j are the imposed equality and inequality constraints and \mathbf{x} denotes the vector of the design variables $(x_1, x_2, x_3, \dots, x_n)^T$. In this study, the design variables \mathbf{x} are:

$$\mathbf{x} = [D, H, W, L, U_{app}, Q]$$

Inequality constraints are:

$$D (mm) \geq 10 \quad (24)$$

$$1.25 \leq \mathcal{S}_L \leq 3 \quad (25)$$

$$1.25 \leq \mathcal{S}_T \leq 3 \quad (26)$$

$$\gamma \geq 20 \quad (27)$$

The objective function can be redefined by using Lagrangian function as follows:

$$\mathcal{L}(\mathbf{x}, \lambda, \chi) = f(\mathbf{x}) + \sum_{j=1}^m \lambda_j g_j(\mathbf{x}) - \sum_{j=m+1}^n \chi_j l_j(\mathbf{x}) \quad (28)$$

where λ_j and χ_j are the Lagrange multipliers. The λ_j can be positive or negative but the χ_j must be ≥ 0 . In addition to Kuhn-Tucker conditions, the other necessary condition for \mathbf{x}^* to be a local minimum of the problem, under consideration, is that the Hessian matrix of \mathcal{L} should be positive semidefinite, i.e.

$$\mathbf{v}^T \nabla^2[(\mathbf{x}^*, \lambda^*, \chi^*)] \mathbf{v} \geq 0 \quad (29)$$

For a local minimum to be a global minimum, all the eigen-values of the Hessian matrix should be ≥ 0 .

A system of non-linear equations is obtained, which can be solved using numerical methods such as a multivariable Newton-Raphson method. This method has been described in Stoecker [24] and applied by Culham and Muzychka [25], Culham et al. [26] and Khan et al. (2004) to study the optimization of plate or pin fin heat sinks. In this study, the same approach is used to optimize the overall performance of a tube bank in such a manner that all relevant design conditions combine to produce the best possible tube bank for the given constraints. The optimized results are then compared for in-line and staggered arrangements.

A simple procedure was coded in MAPLE 9, a symbolic mathematics software, which solves the system of N non-linear equations using the multivariable Newton-Raphson method. Given, the Lagrangian \mathcal{L} , the solution vector $[\mathbf{x}]$, initial guess $[\mathbf{x}_0]$, and maximum number of iterations N_{max} , the procedure systematically applies the Newton-Raphson method until the desired convergence criteria and/or maximum number of iterations is achieved. The method is quite robust provided an adequate initial guess is made.

IV. Results and Discussion

The objective of this study is to determine an optimal tube bank by minimizing the dimensionless entropy generation rate for different design variables including the tube diameter, D , tube length, L , cross-sectional area of tube bank, $W \times H$, and heat load, Q . In each case, optimum approach velocity/Reynolds number is determined corresponding to the minimum entropy generation rate. It is assumed that hot water is passed through the tubes, while air is passed in cross flow over the tubes. The ambient temperature and the tube wall temperatures are fixed at $300K$ and $365K$ respectively. The problem is solved for three different longitudinal and transverse pitch ratios and the overall performance is compared for both the in-line and staggered arrangements.

The dimensions/data given in Table 1 are used as the default case to determine the performance parameters for both in-line and staggered tube banks. In the first case, three different tube diameters 12, 13, and 14 mm are considered with the constraints of a maximum tube bank volume to dissipate a heat load of $20kW$. The problem is solved for each diameter corresponding to three dimensionless pitch ratios 1.25×1.25 , 1.5×1.5 , and 2.0×2.0 . The results of optimization for the dimensionless heat transfer rate, pressure drop and the number of tubes are tabulated in Tables 2 and 3 for each arrangement.

These results show that the number of tubes in a given volume decreases with the increase in tube diameter and/or dimensionless pitch ratio. In in-line arrangement, the heat transfer increases with the

Table 1. Dimensions/Data Used for Optimal Design of Tube Banks

Quantity	Dimension/Data
Cross-Sectional Area (mm^2)	235×235
Length of Tubes (mm)	1000
Tube Diameter (mm)	12
Heat Load (kW)	20
Ambient Temperature (K)	300
Tube Wall Temperature (K)	365

Table 2. Results Of Optimization For In-Line Tube Banks

Dimensionless Pitch Ratio $\mathcal{S}_T \times \mathcal{S}_L$	Tube Diameter (mm)	Optimum Approach Velocity (m/s)	Number of Tubes $N_T \times N_L$	Nu_D	ΔP (Pa)	$N_s \times 10^{10}$
1.25×1.25	12	3.4	15×15	88.4	590.2	0.180
	14	3.8	13×13	100.9	621.6	0.191
	16	4.2	11×11	113.1	650.3	0.201
1.5×1.5	12	5.7	13×13	88.5	480.4	0.251
	14	6.4	11×11	101.0	507.3	0.266
	16	7.0	10×10	112.9	532.8	0.281
2.0×2.0	12	9.6	10×10	91.5	452.7	0.365
	14	10.63	8×8	104.0	477.7	0.389
	16	11.6	7×7	116.0	499.9	0.410

Table 3. Results Of Optimization For Staggered Tube Banks

Dimensionless Pitch Ratio $\mathcal{S}_T \times \mathcal{S}_L$	Tube Diameter (mm)	Optimum Approach Velocity (m/s)	Number of Tubes $N_T \times N_L$	Nu_D	ΔP (Pa)	$N_s \times 10^{10}$
1.25×1.25	12	2.8	15×15	122.2	657.6	0.179
	14	3.2	13×13	142.0	660.8	0.180
	16	3.64	11×11	161.7	664.5	0.181
1.5×1.5	12	5.1	13×13	105.5	535.1	0.254
	14	5.8	11×11	121.8	544.1	0.259
	16	6.6	10×10	137.9	553.1	0.265
2.0×2.0	12	8.43	10×10	90.7	580.3	0.441
	14	9.48	8×8	104.0	597.6	0.458
	16	10.5	7×7	116.8	612.5	0.473

increase in tube diameter and/or dimensionless pitch ratio, whereas in staggered arrangement, the heat transfer increases with the tube diameter but decreases with the increase in dimensionless pitch ratio. In both arrangements, the pressure drop increases with the increase in tube diameter but decreases with the increase in dimensionless pitch ratio. Results of optimization are also shown in Figs. 2(a) and 2(b) for each

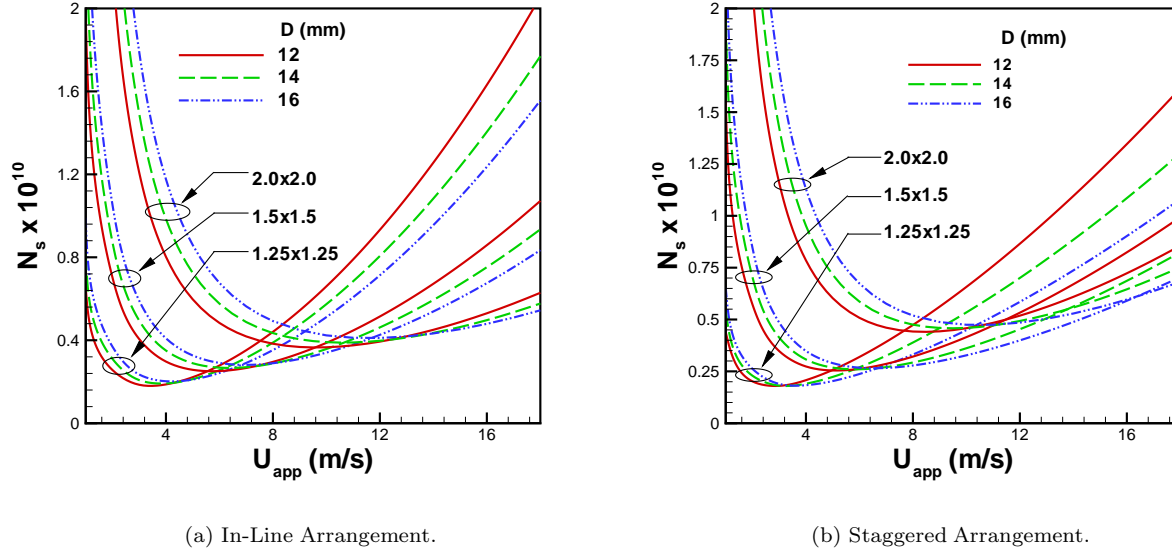
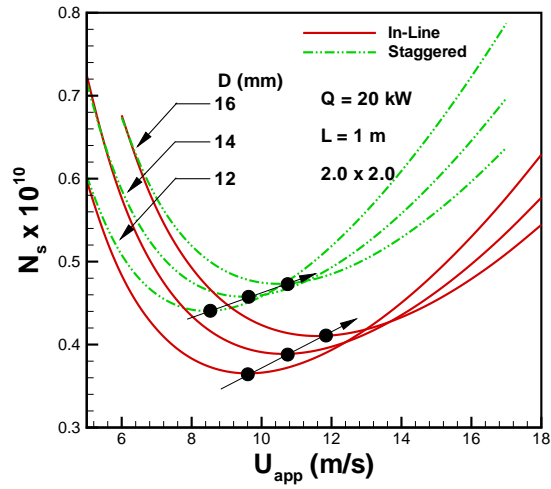


Figure 2. Dimensionless Entropy Generation Rate for Both Arrangements.

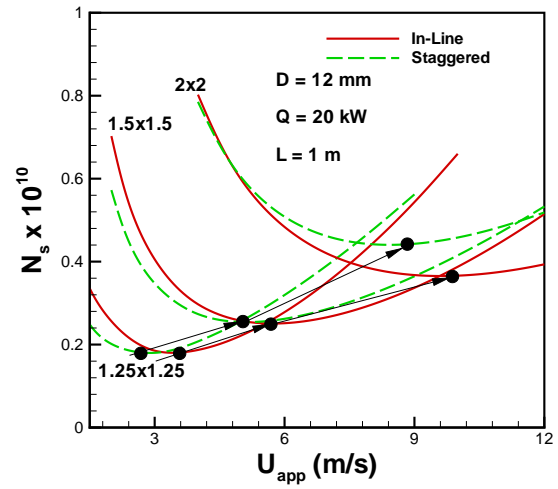
arrangement. It is clear from Fig. 2 that each arrangement has optimum approach velocity for each tube diameter and for each dimensionless pitch ratio. As the tube diameter and/or the dimensionless pitch ratio increases, the optimum approach velocity as well as the dimensionless entropy generation rate increase.

For a given dimensionless pitch ratio 2.0×2.0 , both arrangements are compared in Fig. 3(a) for three tube diameters. It shows that the optimum approach velocity increases with the tube diameter for both arrangements. The dimensionless entropy generation rate is higher for staggered arrangement. In Fig. 3(b), both arrangements are compared for a given tube diameter. It shows that the staggered arrangement gives better performance for lower pitch ratios and lower approach velocities but for higher approach velocities and widely spaced tube banks, in-line arrangement are better with reference to lowest dimensionless entropy generation rate. Optimum approach velocities are also found to be lower for staggered arrangement in case of compact banks. Figures 4(a) and 4(b) show the effects of heat load on the performance of compact and widely spaced tube banks for both arrangements. The optimum dimensionless entropy generation rate decreases with the increase in heat load and the dimensionless pitch ratio for both arrangements. In case of compact tube banks, the staggered arrangement gives better performance for lower approach velocities only, whereas in-line arrangement is better for higher approach velocities and also for widely spaced banks. The optimum approach velocities increases with the heat loads for both type of banks.

The effects of tube length on the performance of tube banks are shown in Figs. 5(a) and 5(b) for both arrangements. Figure 5(a) shows those effects for a compact tube bank. It shows that the effect of tube length is almost negligible on the optimum dimensionless entropy generation rate but the optimum approach velocity decreases with the increase in tube length for both arrangements. In case of widely spaced tube bank, the in-line arrangement performs much better for all three cases. Again, the optimum approach velocity decreases with the increase in tube length for both arrangements. Figure 6 shows the effects of Reynolds number on the performance of tube banks for both arrangements. For different dimensionless

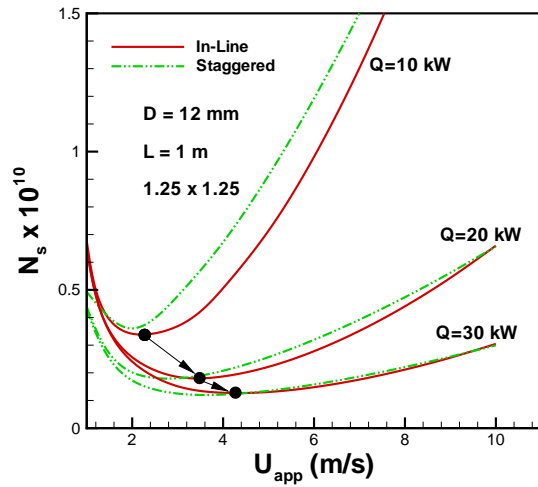


(a) For Given Tube Diameter.

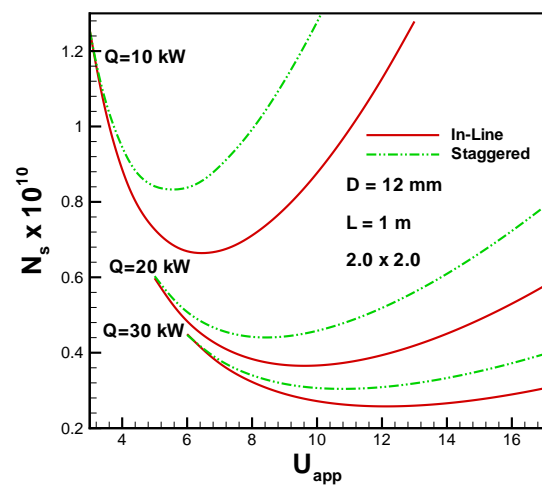


(b) For Given Pitch Ratio.

Figure 3. Comparison of In-Line and Staggered Arrangements for Given Pitch Ratio and Tube Diameter.

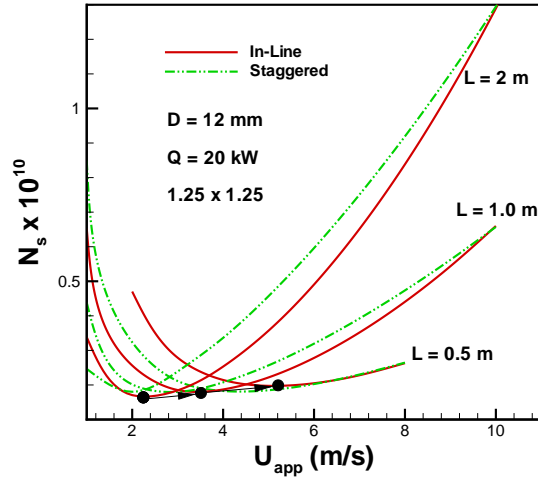


(a) For 1.25×1.25 .

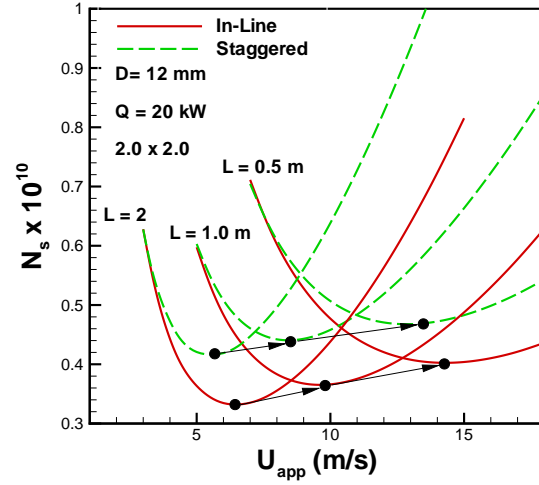


(b) For 2.0×2.0 .

Figure 4. Comparison of In-Line and Staggered Arrangements for Compact and Widely Spaced Tube Banks.



(a) For 1.25×1.25 .



(b) For 2.0×2.0 .

Figure 5. Effect of Tube Length on Dimensionless Entropy generation Rate.

bank heights, in-line arrangement gives better performance for lower Reynolds numbers but as the Reynolds number increases, staggered arrangement does good job for widely spaced tube bank. It shows that the optimum Reynolds number increases with the decrease in the cross-sectional area.

V. Conclusions

A scientific procedure is presented for determining optimal design of tube banks for both the in-line and staggered arrangements. The effects of tube diameter, tube length, dimensionless pitch ratios, front cross-sectional area of the tube bank, and heat load are examined with respect to its role in influencing optimum design conditions and the overall performance of the tube bank. It is demonstrated that the staggered arrangement gives better performance for lower approach velocities and longer tubes, whereas in-line arrangement performs better for higher approach velocities and larger dimensionless pitch ratios. Compact tube banks perform better for both arrangements and for smaller tube diameters.

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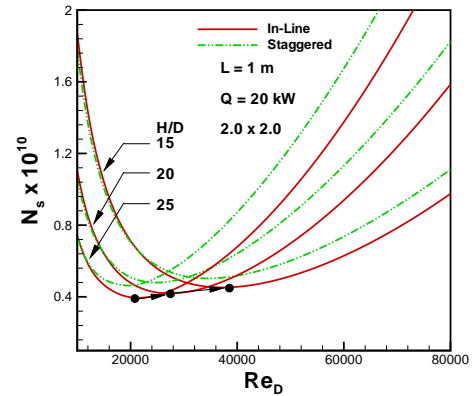


Figure 6. Effect of Reynolds Number on Dimensionless Entropy generation Rate for Widely Spaced Tube Bank.

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